

Quantum
Information
Security

Andreas Winter

(ICREA & Universitat Autònoma de Barcelona)

Quantum
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(kind of a tutorial)

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Outline

0. Quantum formalism

1. Motivation: quantum privacy amplification

2. Tools: (smoothed) min-entropy

3. Applications: QKD, quantum data hiding

0. Quantum formalism

We only need the "probabilistic core" of quantum theory, not the "mechanics".

[See (e.g.) M. M. Wilde, QIT, Cambridge University Press]

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Non-commutative probability theory:

- States (densities)
- Transformations (channels)
- Observable probabilities
(measurements & Born rule)

[See (e.g.) M. M. Wilde, QIT, Cambridge University Press]

0. Quantum formalism

Each quantum system characterized by a complex Hilbert space \mathcal{H} , A , B , C , ...

Here: dimensions $|\mathcal{H}|$, $|A|$, ... $< \infty$.

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Dirac ("bra-ket") notation:

- Vectors $|\psi\rangle$

- Covectors $\langle\varphi|$

- Inner product $\langle\varphi|\psi\rangle$

- Outer product $|\psi\rangle\langle\varphi|$, in particular
rank-one projector $|\psi\rangle\langle\psi|$

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- Hermitian conjugate \dagger :

$$|\psi\rangle^\dagger = \langle\psi|, \quad \langle\varphi|^\dagger = |\varphi\rangle$$

- Everything described in terms of operators $A \in L(\mathcal{H})$, $A^\dagger \in L(\mathcal{H})$.

[See (e.g.) M. M. Wilde, QIT, Cambridge University Press]

0-a. States

States on \mathcal{H} are *density operators*:

$\rho \geq 0$ (positive semidefinite), $\text{Tr } \rho = 1$.

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(Well, it's really matrices...)

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Extremal points: rank-one projectors

$|\psi\rangle\langle\psi|$ - *pure states*.

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With a distinguished ONB $|x\rangle$ ("computational basis"), have diagonal matrices $\rho = \sum_x p_x |x\rangle\langle x|$, isomorphic to the probability simplex.

[See (e.g.) M. M. Wilde, QIT, Cambridge University Press]

0-6. Measurement

Observations on quantum systems require an **observable**, which is a resolution of the identity (aka POVM):

$$M_y \geq 0, \text{ with } \sum_y M_y = I.$$

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State as the "catalog of expectations"
(Schrödinger)

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Example: Discriminating equiprobable hypotheses ρ and σ . Need a binary POVM $(M, 1-M)$, with $0 \leq M \leq 1$.

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$$= \frac{1}{2} + \frac{1}{4} \underbrace{\text{Tr} (\rho - \sigma)(2M - 1)}$$

$$\max = \|\rho - \sigma\|_1 \quad (\text{trace norm})$$

$$= \text{Tr} |\rho - \sigma|$$

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O-c. Composite systems

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The latter allows us to identify a POVM (M_y) on A with $(M_y \otimes I)$ on $A \otimes B$. And for (N_z) on B , joint POVM $(M_y \otimes N_z)$.

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Where does that leave us with the states? $S(A \otimes B)$ clearly contains all tensor product densities $\alpha \otimes \beta$ and their convex hull, but it has so much more...

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Where does that leave us with the states? $S(A \otimes B)$ clearly contains all tensor product densities $\alpha \otimes \beta$ and their convex hull, but it has so much more...

Entanglement! For instance all pure states $|\psi\rangle\langle\psi|$ with $|\psi\rangle \neq |\alpha\rangle \otimes |\beta\rangle$.

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Composition of A and B (Hilbert spaces) is a system with tensor product Hilbert space $A \otimes B$.

Marginals: For a state $\rho^{AB} \in S(A \otimes B)$, the rule

$$\text{Tr} \rho(X \otimes I) = \text{Tr} \rho^A X \quad (\text{for all } X)$$

singles out uniquely the partial trace:

$$\rho^A = \text{Tr}_B \rho^{AB}.$$

[See (e.g.) M. M. Wilde, QIT, Cambridge University Press]

O-c. Composite systems

Purification: For every state $\rho \in S(A)$, there exists a unit vector $|\varphi\rangle \in A \otimes A'$, s.t. ρ is the marginal of the pure state:

$$\rho^A = \text{Tr}_{A'} |\varphi\rangle\langle\varphi|.$$

Any other pure state $|\psi\rangle\langle\psi|$ with this property is related to $|\varphi\rangle\langle\varphi|$ by a unitary conjugation on A' :

$$|\psi\rangle = (1 \otimes U) |\varphi\rangle.$$

[See (e.g.) M. M. Wilde, QIT, Cambridge University Press]

0-d. Quantum channels

Like classical channels, quantum channels map densities to densities, i.e.

$$N: S(A) \rightarrow S(B).$$

As it must respect convex combinations, extends uniquely a linear, positive, trace preserving map $N: L(A) \rightarrow L(B)$.

However, want complete positivity, i.e.

$$N \otimes \text{id}_C: S(A \otimes C) \rightarrow S(B \otimes C)$$

should be positive for all auxiliary C .

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Quantum channel

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is a linear, completely positive and trace preserving (cptp) map.

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Basic example: unitary conjugation on A

$$\mathcal{N}(\rho) = U \rho U^\dagger.$$

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Basic example: unitary conjugation on A

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Another one: partial trace from AB to A

$$N(\rho) = \text{Tr}_B \rho.$$

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0-d. Quantum channels

More examples:

1) Noiseless channel = identity id_A .

2) Constant channel $K(\rho) = \omega_0$.

3) Depolarizing channels

4) Amplitude damping channels

5) Phase damping channels

6) Erasure channel $\mathcal{E}_g(\rho) = (1-g)\rho \oplus g|* \rangle \langle *|$

[See (e.g.) M. M. Wilde, QIT, Cambridge University Press]

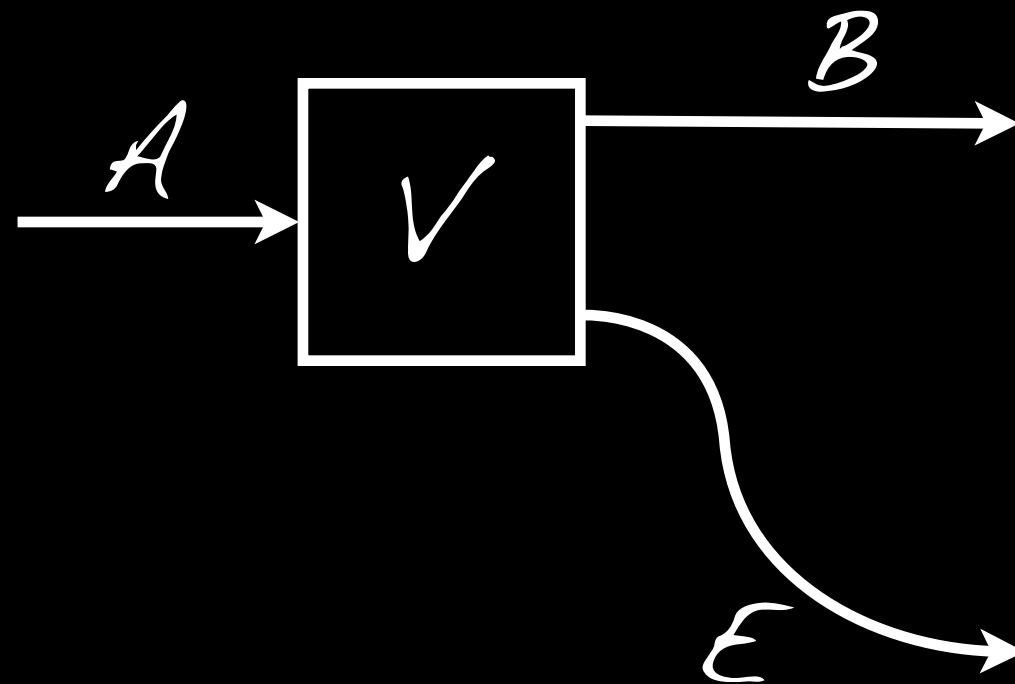
0-d. Quantum channels

Stinespring: $N(\rho) = \text{Tr}_E V \rho V^\dagger$,
with an isometry $V: A \hookrightarrow B \otimes E$.

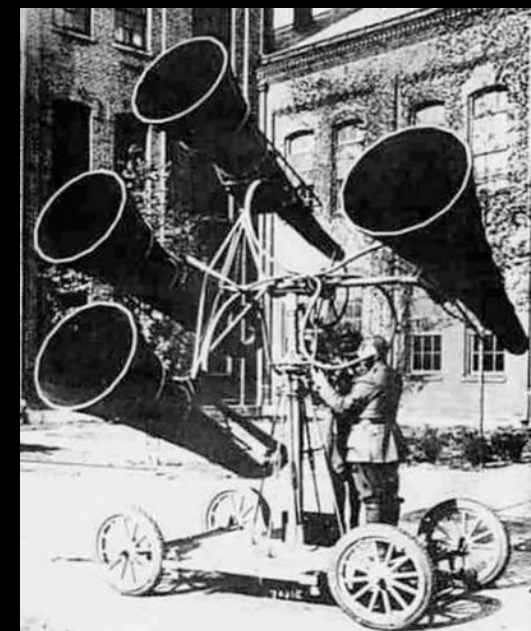
Complementary channel:

$$\hat{N}(\rho) = \text{Tr}_B V \rho V^\dagger.$$

Alice

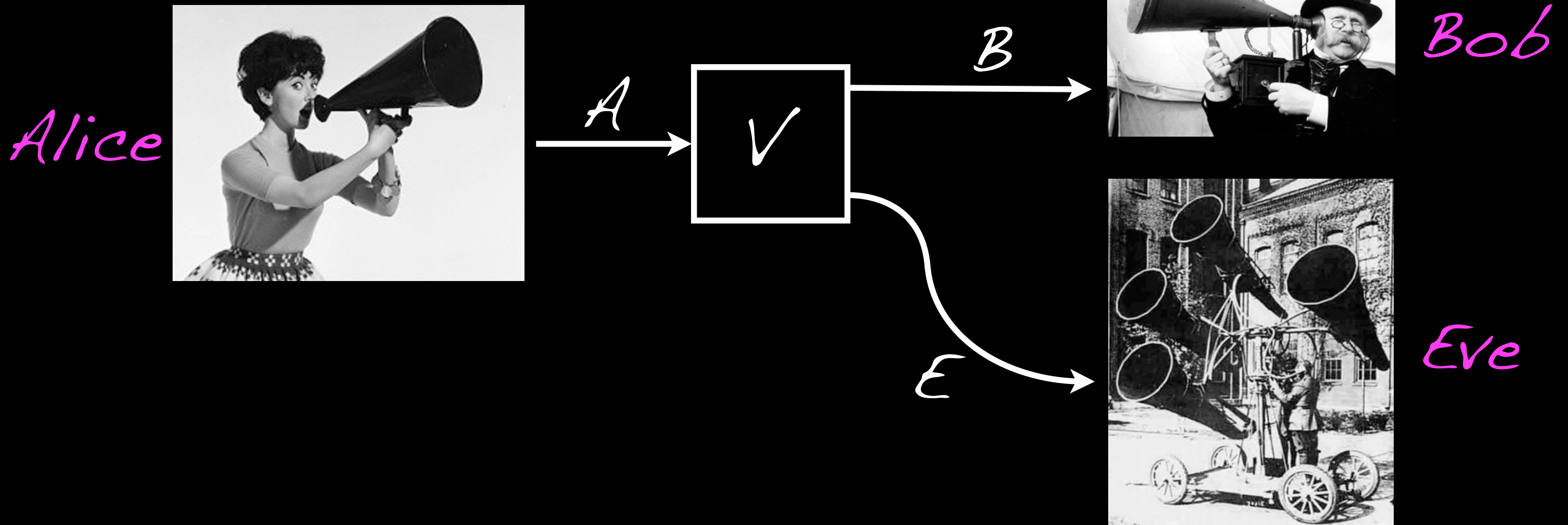


Bob



Eve

0-d. Quantum channels



It's a wiretap model! Note that automatically if Bob's channel is (nearly) noiseless, Eve's is (nearly) constant.

0-d. Quantum channels

Important property of cptp maps: trace norm is contractive. I.e., for any channel N and any operator ξ ,

$$\|N(\xi)\|_1 \leq \|\xi\|_1.$$

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Important property of cptp maps: trace norm is contractive. I.e., for any channel N and any operator ξ ,

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Other metrics with this property, e.g. the purified distance

$$P(\rho, \sigma) = \min \frac{1}{2} \|\varphi - \psi\|_1 \text{ s.t. purifications } \varphi, \psi \text{ of } \rho, \sigma$$

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1. Privacy amplification

Alice has a "raw key" X (maybe shared with Bob), which is compromised: conditional on side information of Eve, the distribution of X is not uniform.

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But if we know $-\log \Pr\{X=x\} \geq h$, then 2-universal hash-function F to k bits yields $Y=F(X)$ with

$$\|P_{YF} - U \otimes P_F\|_1 \leq 2^{-(h-k)/2},$$

where U is the uniform distribution.

1. Privacy amplification

Including Eve explicitly, have Alice's X and Eve's Z jointly distributed. The optimal h is

$$H_{\min}(X|Z) = -\log \max_{x,z} \Pr\{X=x|Z=z\}$$

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What can we do if Eve has quantum information? What can that conditional probability mean?

1. Privacy amplification

Now Alice and Eve have a joint *state*:

$$\rho^{XE} = \sum_x p_x |x\rangle\langle x|^X \otimes p_x^E.$$

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It turns out the right definition is about comparing with operators $\lambda(1 \otimes \sigma)$:

$$H_{\min}(X|E) = -\log \min_{\lambda, \sigma} \lambda \text{ s.t. } \rho \leq \lambda(1 \otimes \sigma),$$

where σ is a state on E and λ is real.

1. Privacy amplification

Let $H_{\min}(X|E) = -\log \min_{\lambda, \sigma} \lambda$ s.t. $\rho \leq \lambda(I \otimes \sigma) \geq h$.

Then 2-universal hash-function F to k bits yields $Y=F(X)$ with

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...essentially optimal without knowing more about the joint state.

[Cf. R. Renner, PhD thesis, arXiv:quant-ph/0512258]

2. Min-entropy calculus

Renner found that one could develop a meaningful one-shot information calculus based on *conditional min-entropy*:

$$H_{\min}(A|B)_\rho = -\log \min_{\lambda, \sigma} \lambda \text{ s.t. } \rho^{AB} \leq \lambda(I^A \otimes \sigma^B),$$

for an arbitrary state ρ on $A \otimes B$.

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for an arbitrary state ρ on $A \otimes B$. Not only is this a semidefinite programme (SDP), as we shall see it shares many good properties with the von Neumann entropy.

[R. Renner, PhD thesis, arXiv:quant-ph/0512258;
M. Tomamichel, PhD thesis, arXiv:1203.2142]

[Entropy digression:

The von Neumann entropy of a state ρ is $S(\rho) = -\text{Tr} \rho \log \rho = H(\vec{\lambda})$, where $\vec{\lambda}$ is the eigenvalue vector of ρ .

We define information expressions by linear-combining entropies, à la Shannon:
E.g. $S(A|B)_\rho = S(\rho^{AB}) - S(\rho^B)$, or conditional mutual information

$$I(A:B|C)_\rho = S(\rho^{AB}) + S(\rho^{BC}) - S(\rho^{ABC}) - S(\rho^C).$$

[See (e.g.) M. M. Wilde, QIT, Cambridge University Press]

[Entropy digression:

While the quantum conditional entropy $S(A|B)_\rho = S(\rho^{AB}) - S(\rho^B)$ can be negative,

the conditional mutual information

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is always ≥ 0 ; this is Lieb-Ruskai's strong subadditivity...

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2. Min-entropy calculus

Conditional (smooth) min-entropy:

$$H_{\min}^{\epsilon}(A|B)_{\rho} = -\log \min_{\lambda, \sigma} \text{Tr } S \text{ s.t. } \rho^{AB} \leq \lambda \otimes S^B$$

In the applications (privacy amplification, ...) we can gain by "smoothing":

$$H_{\min}^{\epsilon}(A|B)_{\rho} = \max H_{\min}(A|B)_{\omega} \text{ s.t. } P(\rho, \omega) \leq \epsilon$$

Amazingly, also this is an SDP! Useful due to SDP duality (...)

[R. Renner, PhD thesis, arXiv:quant-ph/0512258;
M. Tomamichel, PhD thesis, arXiv:1203.2142]

2. Min-entropy calculus

Motivated by the von Neumann entropy identity $S(A|B) = -S(A|C)$ for a pure state on ABC , we define also a conditional (smooth) max-entropy:

$$H_{\max}^{\epsilon}(A|B)_{\rho} = -H_{\min}^{\epsilon}(A|C)_{\varphi} \text{ w.r.t. purification } \varphi$$

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As in the classical case, "min" applies to extraction problems, while "max" to compression problems...

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2. Min-entropy calculus

These quantities share many features with the von Neumann entropy, e.g. strong subadditivity reads $S(ABC) \leq S(AB)$:

$$H_{\min}^{\epsilon}(ABC)_{\rho} \leq H_{\min}^{\epsilon}(AB)_{\rho};$$

$$H_{\max}^{\epsilon}(ABC)_{\rho} \leq H_{\max}^{\epsilon}(AB)_{\rho}.$$

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2. Min-entropy calculus

$$\mathcal{H}_{\min}^{\varepsilon}(A|BC)_{\rho} \leq \mathcal{H}_{\min}^{\varepsilon}(A|B)_{\rho};$$

$$\mathcal{H}_{\max}^{\varepsilon}(A|BC)_{\rho} \leq \mathcal{H}_{\max}^{\varepsilon}(A|B)_{\rho}.$$

Furthermore, while $\mathcal{H}_{\min}^{\varepsilon}$ increases and $\mathcal{H}_{\max}^{\varepsilon}$ decreases with ε , we have for suff. small ε and δ ,

$$\mathcal{H}_{\min}^{\varepsilon}(A|B)_{\rho} \leq \mathcal{H}_{\max}^{\delta}(A|B)_{\rho} + O(1).$$

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2. Min-entropy calculus

Using properties of typical subspaces, one can also show the "quantum AEP": For n -tensor power states $\rho^{\otimes n}$ on $A^n B^n$, and $0 < \epsilon < 1$,

$$\lim_{n \rightarrow \infty} \frac{1}{n} \mathcal{H}_{\min}^{\epsilon}(A^n | B^n) = S(A|B)_{\rho},$$

$$\lim_{n \rightarrow \infty} \frac{1}{n} \mathcal{H}_{\max}^{\epsilon}(A^n | B^n) = S(A|B)_{\rho}.$$

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3. Applications: QKD, ...

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Since then it was used in condensed matter physics [Area laws, Brandão/Horodecki] and for strong converses [talk by Ciara Morgan].

3'. LOCC hiding efficiency

Recall the state discrimination problem, but now let ρ and σ be states on a composite system: $\mathcal{H} = A \otimes B$, and let A (Alice) and B (Bob) be far apart.

Without a quantum channel between them, they can only perform **local operations and classical communication (LOCC)**.

[Walgate et al., PRL85:4972-4975, 2000]

[Terhal/DiVincenzo/Leung, PRL86:5807-5810, 2001]

3'. LOCC hiding efficiency

ρ or σ



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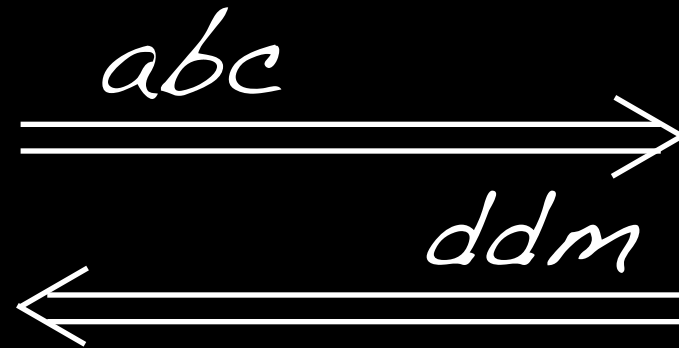


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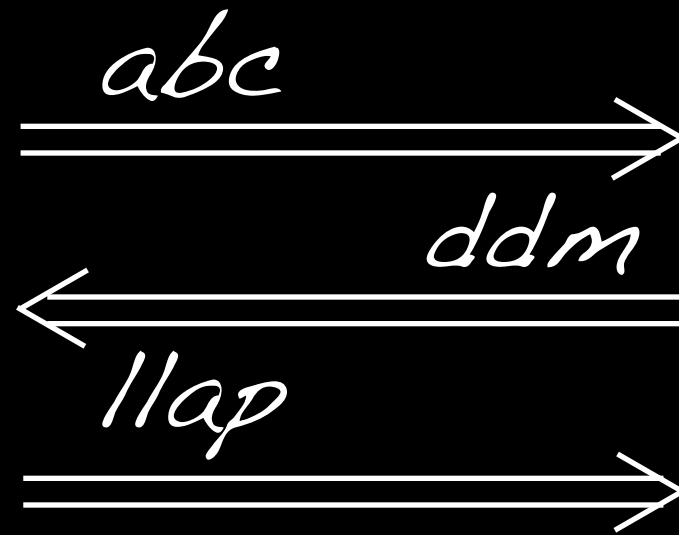


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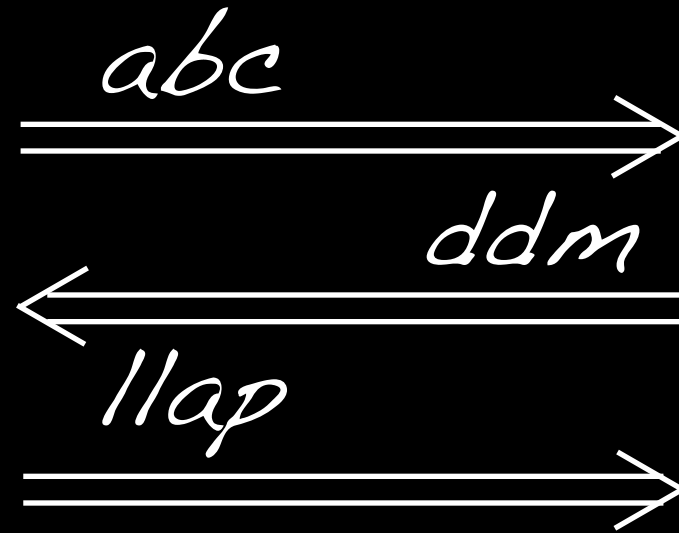


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$$\max P\{\text{guess}\} =: \frac{1}{2} + \frac{1}{4} \|\rho - \sigma\|_{\text{LOCC}}$$

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3'. LOCC hiding efficiency

ρ or σ



↓
guess

$$\max P\{\text{guess}\} =: \frac{1}{2} + \frac{1}{4} \|\rho - \sigma\|_{\text{LOCC}}$$

Can be much smaller than $\|\cdot\|_1$: data hiding!

[Terhal/DiVincenzo/Leung, PRL86:5807-5810, 2001]

3'. LOCC hiding efficiency

Let ρ_x be states on a composite system $\mathcal{H} = A \otimes B$, with the systems of Alice and Bob having dimensions $|A| = |B| = d$ ($x = 1 \dots M$), such that any pair of states is LOCC-hiding, but that there is a global POVM revealing x among all M .

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How large can M be in relation to d ?

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Can we have more? **NO!**

[AW, in preparation, 2014-2015]

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Here "(2,2)" - note that LOCC data hiding cannot be perfect (re secrecy).

Formal statement: If $\epsilon, \delta > 0$ are small enough, s.t. $\|\rho_x - \rho_y\|_{LOCC} \leq \delta$ for all x, y , and $\text{Tr} \rho_x \mathcal{D}_x \geq 1 - \epsilon$ for a POVM (\mathcal{D}_x) .

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Then: $\mu, C > 0$ (dep. only on ϵ, δ) with $\log(M-1) \leq \mathcal{H}_{\max}^{\mu}(A|B)_{\Omega} + C \leq \log d + O(1)$, w.r.t. uniform average state Ω of the ρ_x .

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$$H_{\min}^{\varepsilon+2\varepsilon'+\varepsilon''}(AB|C)_\rho \geq H_{\min}^{\varepsilon'}(A|BC)_\rho + H_{\min}^{\varepsilon''}(B|C)_\rho - \log \frac{2}{\varepsilon^2}, \quad (5.11)$$

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Proof (sketch):

Define $\Omega^{XAB} = \frac{1}{M} \sum_x |x\rangle\langle x|^X \otimes \rho_x^{AB}$. Then:

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{Decodability}

QED

Compare secret sharing [A. Shamir, 1979; G.R. Blakley, 1979]: In any tight secret sharing scheme (i.e. where all subsets are either authorized or adversarial), each relevant share must be at least as large as the secret.

Corollary: In any tight LOCC data hiding scheme, each relevant share must be at least as large as the secret (plus a constant). Attained for " (n,t) " scheme.

[AW, in prep.; Hayden/Leung/Smith, PRA 71:062339, 2005]



Thank you for your attention

